# <span id="page-0-0"></span>Richard Feynman, the Path Integral, and Least Action **Principles**

#### Ross Dempsey

Department of Physics and Astronomy Johns Hopkins University

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## **Outline**

**1** [Least Action Principles](#page-2-0)

2 [Quantum Mechanics](#page-16-0)

3 [Path Integral Formulation](#page-27-0)

- <span id="page-2-0"></span> $\bullet$  Newton discovered that  $\boldsymbol{F} = m \boldsymbol{a}$ . This tells us how particles will move at each point in time
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- What governs the overall trajectories?
- Newton's laws are equivalent to minimizing a quantity, *action*, which depends on whole trajectories
- No point in time is more special than the others, so the action should care equally about every time slice
- Write the action as an integral (sum) over all times:

$$
S=\int \mathcal{L}\, dt
$$

 $\circ$   $\mathcal{L}$  is the Lagrangian

#### Lagrangians and Theories

Mathematics gives a recipe for converting Lagrangians to equations of motion:

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\frac{d}{dt}\frac{\partial \mathcal{L}}{\partial \dot{x}} = \frac{\partial \mathcal{L}}{\partial x}.
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- $\bullet$  Example: in mechanics,  $\mathcal{L} = T V$ , the difference of kinetic and potential energies
- Resulting equations of motion:

$$
m\ddot{x} = -\frac{\partial V}{\partial x} = F
$$

Newton's second law!

#### Electromagnetism

- Sometimes this is called "Lagrangian mechanics." This obscures a crucial fact: all known physical theories can be approached in the same way
- Example: Lagrangian for electromagnetism is

$$
\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + 4\pi j_{\mu} A^{\mu} = \frac{1}{2} (E^2 - B^2)
$$

Using the recipe, this gives the dynamic Maxwell equations in a vacuum:

$$
\nabla \cdot \mathbf{E} = 4\pi \rho, \qquad \nabla \times \mathbf{B} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}
$$



Figure: Maxwell's equations on Etsy

## Why care?

- If we already know Newton's laws and Maxwell's laws, why generate them with a Lagrangian?
- Lagrangians can be extremely powerful for understanding conservation laws
- Look again at the equations of motion:

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- Independence of  $x$  generates a conserved quantity  $\frac{\partial \mathcal{L}}{\partial \dot{x}}$
- This is *Noether's theorem*: a symmetry generates a conservation law
- By applying Noether's theorem to time translation symmetry, we obtain conservation of energy
- The energy is given by

$$
\mathcal{H} = \sum \frac{\partial \mathcal{L}}{\partial \dot{x}_i} \dot{x}_i - \mathcal{L}
$$

**This is also called the Hamiltonian** 

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- $\bullet$  "The president says, 'That's a very good idea. On the fraternity word of honor!' So he goes around the table and asks each guy, one by one:"
- "Feynman, did you take the door?" "Yeah, I took the door." "Cut it out, Feynman; this is serious!"



## <span id="page-16-0"></span>Quantum States

- Classical states have definite values for every observable
- Quantum theory allows for superpositions of definite states
- When expressed in terms of position states, quantum states are called wavefunctions



#### Example: Hydrogen



- Superpositions are interpreted as specifying probabilities
- Measurement of a quantum system forces it to collapse into a definite state, based on the probabilities
- Until it is measured, the system evolves as a superposition

## Time Evolution

○ Quantum systems evolve according to the Schrödinger equation:

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i\hbar\frac{\partial}{\partial t}\ket{\psi}=\mathcal{H}\ket{\psi}
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## Time Evolution

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- $\circ$  The operator  ${\mathcal H}$  is the quantum version of the Hamiltonian
- o If a system is in a definite energy state, it evolves in a trivial way
- In the most general case, time evolution can be expressed in terms of an operator  $U$ :

$$
|\psi(t_1)\rangle = U(t_1,t_0) |\psi(t_0)\rangle
$$

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$$
\mathcal{L} = \frac{i\hbar}{2} \left( \psi^* \dot{\psi} - \dot{\psi}^* \psi \right) - \frac{\hbar^2}{2m} \nabla \psi^* \nabla \psi - V(\mathbf{r}, t) \psi^* \psi
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- **•** "Beauty is the first test." This fails.

## Break: Feynman and the Safe

- Feynman worked on the Manhattan project in Los Alomos during World War II
- "To demonstrate that the locks meant nothing, whenever I wanted somebody's report and they weren't around, I'd just go in their office, open the filing cabinet, and take it out. When I was finished I would give it back to the guy: "Thanks for your report."'
- "I opened the safes which contained all the secrets to the atomic bomb: the schedules for the production of the plutonium, the purification procedures, how much material is needed, how the bomb works..."



## <span id="page-27-0"></span>Double Slit Experiment

Since quantum mechanics allows for superpositions, an electron can go through two slits at once to get to a screen on the opposite side



Before we wrote time evolution in a single step:

$$
|\psi(t_1)\rangle = U(t_1,t_0) |\psi(t_0)\rangle
$$

 $\bullet$  We could also break up the time evolution into  $N$  steps:

 $|\psi(t_N)\rangle = U(t_N, t_{N-1})U(t_{N-1}, t_{N-2})\cdots U(t_1, t_0)|\psi(t_0)\rangle$ 

## Slicing up the Double Slit

- Breaking up time evolution is somewhat like making a measurement at multiple points along a path
- One way of making a measurement is to use a double slit apparatus and check which slit the particle goes through



## Turning off Measurements

- If we don't make measurements at the slits, the particle can be in a superposition
- The particle can take all possible paths through the slits



## Fine-Grained Slits

- $\circ$  Instead of a double slit, how about three? Four?  $N$ ?
- By adding slits, we increase the freedom of the particle to take multiple paths
- **In the limit, there are no slits, just an open space**



### Summing over Paths

- We now have a new picture of quantum mechanics: particles taking all possible paths from point  $A$  to point  $B$
- If this is how particles work, how come classically we see a single path?



## Classical Limit

In the sum over paths, each one is weighted by a *phase factor*. Formally,

$$
\langle f|i\rangle = \int \mathcal{D}x(t)e^{iS[x(t)]}
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- $\bullet$  Except at the path where S is minimized



## <span id="page-35-0"></span>Classical Limit

- $\bullet$  The quantity  $S$  is the *classical action for the path*
- The path integral reproduces the principle of least action

